

Measurement of antiproton annihilation on Cu, Ag and Au with emulsion films

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ABSTRACT: The characteristics of low energy antiproton annihilations on nuclei (e.g. hadronization and product multiplicities) are not well known, and Monte Carlo simulation packages that use different models provide different descriptions of the annihilation events. In this study, we measured the particle multiplicities resulting from antiproton annihilations on nuclei. The results were compared with predictions obtained using different models in the simulation tools GEANT4 and FLUKA. For this study, we exposed thin targets (Cu, Ag and Au) to a very low energy antiproton beam from CERN's Antiproton Decelerator, exploiting the secondary beamline available in the AEGIS experimental zone. The antiproton annihilation products were detected using emulsion films developed at the Laboratory of High Energy Physics in Bern, where they were analysed at the automatic microscope facility. The fragment multiplicity measured in this study is in good agreement with results obtained with FLUKA simulations for both minimally and heavily ionizing particles.

KEYWORDS: Particle detectors; Emulsion detectors; Antiproton annihilations.

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1. Introduction

Emulsion films have recently been considered as possible position detectors for low-energy anti-matter studies. These studies include the AEgIS (AD6) experiment at CERN [1, 2, 3, 4], whose goal is the measurement of the Earth's gravitational acceleration on antihydrogen atoms. Another collaboration proposed emulsions for their studies on positrons, as described in [5]. In particular, in the case of the AEgIS experiment, the position-sensitive detector must have a micrometer-level resolution to allow the required sensitivity of $\sim 1\%$ for the gravitational acceleration measurement. Spatial resolutions of $\sim 1\text{-}2\ \mu\text{m}$ can be achieved with emulsion films [6], and they have been exploited before for the reconstruction of antihydrogen impact points from annihilation products [3]. Films with this resolution, combined with a time of flight detector, could allow the experimental goal to be achieved. In the same paper, a preliminary study of antiproton-nuclei annihilations was also reported. That study assessed particle multiplicities resulting from antiproton annihilations on emulsion films and aluminium. Recently, again within the context of the AEgIS experiment, similar measurements were performed by means of a silicon detector also acting as an annihilation target [7]. Apart from the obvious applications in nuclear physics, measuring the decay products of low-energy antiproton annihilation in different materials provides a useful check of the ability of standard Monte Carlo packages to reproduce fragment multiplicities, type and energy distributions stemming from antiproton (or antineutron) annihilations on nuclei at rest. Although measurements of the multiplicities of pions and other charged particles with energies higher than $\sim 50\ \text{MeV}$ are available in the literature [8, 9], the production of highly ionizing nuclear fragments with short range has not been studied sufficiently. A measurement of the multiplicities of charged products in antiproton-aluminium annihilations was reported in [3], although only 43 events were analysed in this study, and the tracking efficiency of the detector was limited to 80%. In this paper, we present the results of a study of the multiplicities of charged annihilation products on different target materials, namely copper, silver and gold, using emulsion detectors at the Antiproton Decelerator (AD [10]) at CERN.

2. Experimental setup

Emulsion detectors were used to study the antiproton annihilation products generated in different materials. Before reaching the targets, the 5.3 MeV antiprotons from the AD ($3 \times 10^7 \bar{p}$ /shot every 100 s) were slowed down using several different titanium and aluminium foils with variable thicknesses. Finally, the beam was collimated in a vacuum test chamber after crossing a titanium vacuum separation window with a thickness of $12 \mu\text{m}$. The emulsion detector was situated at the downstream end of the vacuum chamber ($\sim 1 \text{ m}$ in length), where it could be reached by a defocused beam of low-energy antiprotons ($\sim 100 \text{ keV}$). This distance from the degrading layer was necessary to reduce the background due to annihilations taking place at the moderator. A sketch of the experimental assembly is shown in Figure 1. The emulsion detector was operated under ordinary vacuum conditions ($10^{-5} - 10^{-6} \text{ mbar}$). The antiproton intensity measured by the detector was approximately $150/\text{cm}^2$ per shot.

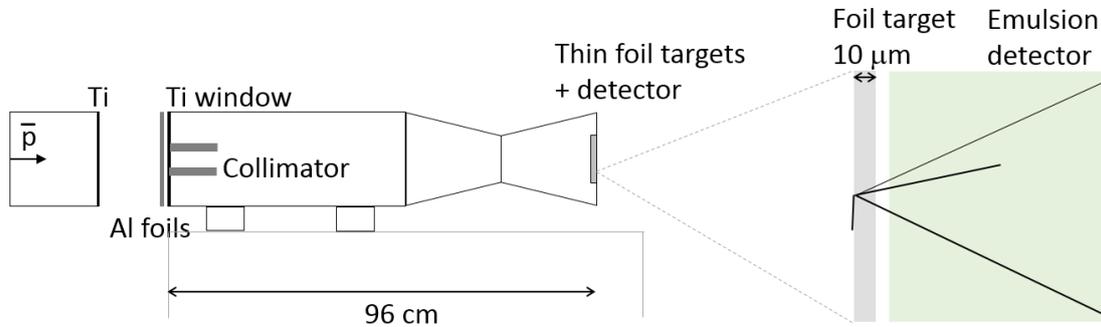


Figure 1. Schematic setup of the experiment. An enlarged view of the target region is shown on the right.

For this study emulsion detectors were produced at the Laboratory for High Energy Physics (LHEP) of the University of Bern by pouring the emulsion gel with a thickness of $\sim 100 \mu\text{m}$, provided by Nagoya University (Japan), on a glass plate (for a review on the emulsion technology see [6]). Glycerin was added to so that the emulsion could operate in vacuum [3]. This emulsion features a very low background with approximately 1-2 thermally induced grains per $1000 \mu\text{m}^3$ [5].

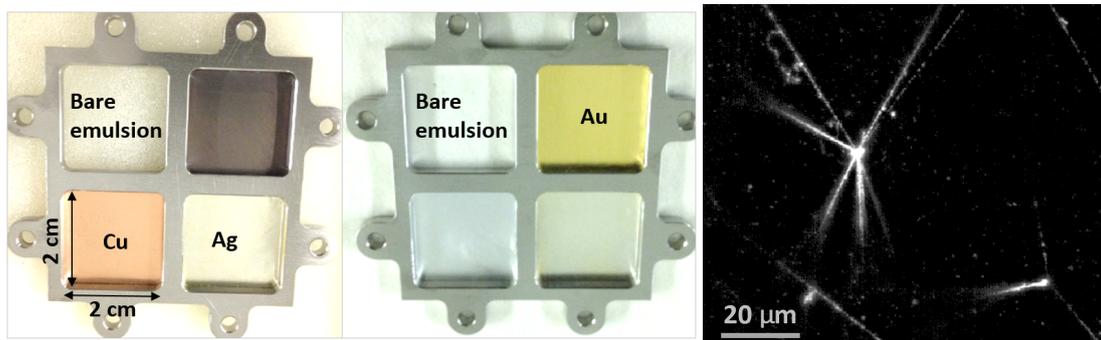


Figure 2. Left and middle: Target arrangement in the two assemblies fixed to the emulsion films. Right: Antiproton annihilations on the bare emulsion surface.

Foils of copper, silver and gold, each having a thickness of $10\ \mu\text{m}$, were placed as targets at the end of the vacuum chamber, in front of the emulsion detectors. Figure 2 shows the targets ($2 \times 2\ \text{cm}^2$ each) fixed to the emulsion film and an example of antiproton annihilating on the emulsion surface. At the antiproton energies obtained after degrading, all annihilations are expected to take place within a few μm of the target surface. During data taking, we collected approximately 1500 antiprotons per cm^2 in about 10 AD shots.

3. Data analysis and results

Data recorded by the emulsion detectors were scanned by an automatic optical microscope and then analysed by exploiting a recently developed fast 3D tracking algorithm [11]. The measured tracking efficiency of our detector was approximately 99% for minimally ionizing particles over a wide angular range, as reported in [11].

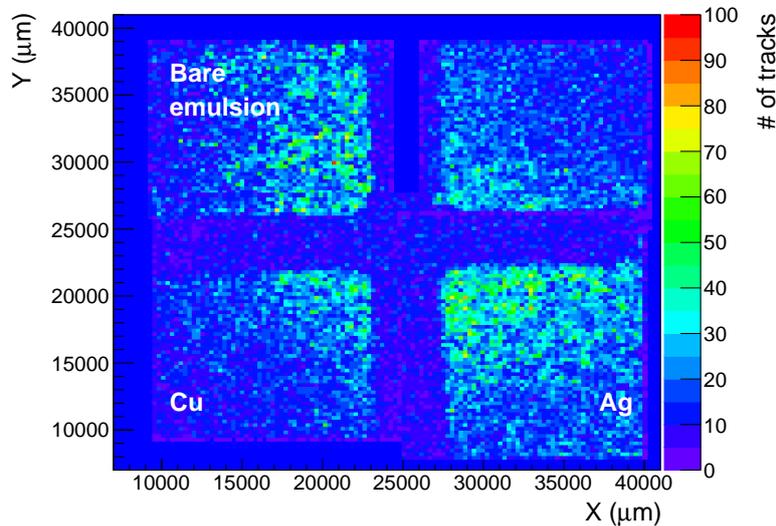


Figure 3. Profile of detected tracks (xy positions of tracks at the emulsion layer) in one of the assemblies. The top right part was not included in our study.

Figure 3 shows the profile of detected tracks (xy positions of tracks at the emulsion layer) in one of the assemblies. Among the reconstructed tracks we only considered those that were longer than $30\ \mu\text{m}$ to avoid considering tracks that were due to the background. An angular cut of $0.4 < \tan\theta < 2.0$ ($22^\circ < \theta < 63^\circ$) was applied, where θ is the track angle with respect to the beam direction. To reconstruct a vertex, at least two three-dimensional tracks were required. The efficiency of vertex reconstruction was estimated by applying the criteria given above to the output of the FLUKA simulation. It was found to be 22% for copper, 24% for silver, and 18% for gold.

Figure 4 (left) shows the distribution of the reconstructed vertex position perpendicular to the film surface (z -direction) in a subarea of the copper target. The peak in z is a measure of the target foil position (the gap between the emulsion surface and the target foil). This measurement was performed by segmenting the analysed area into smaller areas since the target foils were neither flat nor in contact with the film surface. The surface topography obtained from the reconstructed

vertices is shown in Figure 5 for copper, silver and gold targets. In our analysis, we only considered regions of the emulsion film surface with the z value smaller than $100 \mu\text{m}$ because the vertex reconstruction efficiency was uniform within a few percent in this region. The analysed fiducial area was 1.68 cm^2 for the copper target, 1.96 cm^2 for silver and 0.80 cm^2 for gold. The fraction of signal vertices became dominant by requiring vertex reconstruction at the position of the target foil.

The nearly flat distribution in Figure 4 (right) is due to combinatorial background in the vertex reconstruction. The main source of the background was due to accidental combinations of tracks coming from annihilations taking place upstream in the apparatus, which were not completely excluded by the angular cut due to the broad angular distribution. The fraction of tracks from signal vertices to all detected tracks was estimated using the vertex finding efficiency described above and found to be 9% for copper, 9% for silver, and 6% for gold. The number of background vertices was estimated using all the detected tracks in the analysed area by subtracting the above signal track fraction, randomizing positions and slopes of the remaining tracks, reconstructing the vertices and counting the number of vertices that mimicked annihilations in the target.

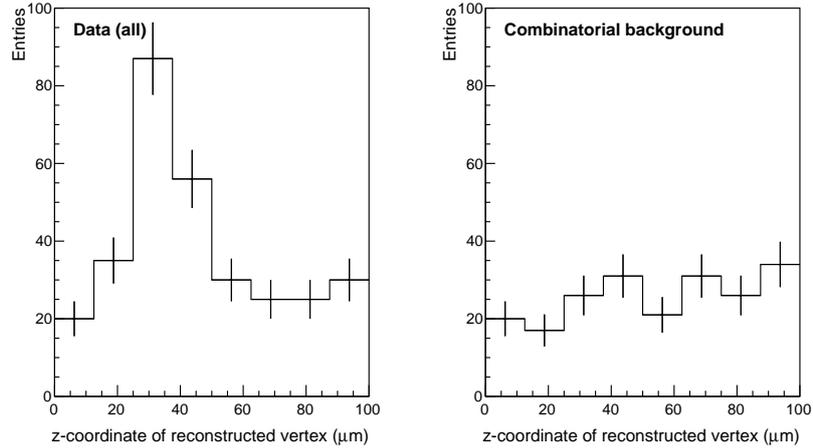


Figure 4. Left: Distribution of the reconstructed vertex position perpendicular to the film surface (z -direction). Right: Estimated combinatorial background.

The background estimated from measured data, which depends on the number of tracks in the event, is shown in the left panes of Figure 6, while the right panes show the multiplicity distributions after subtraction of the background, compared with the Monte Carlo predictions based on the CHIPS [12] and FTFP (FTFP_BERT_TRV) [13] models in the GEANT4 (4.9.5.p02) and FLUKA (2011.2c) [14] frameworks. A total of 654 signal annihilation vertices were reconstructed for copper, 941 for silver and 233 for gold.

We were also able to discriminate between heavily ionizing particles (HIPs) such as protons and nuclear fragments and minimally ionizing particles (MIPs), namely pions. Continuous dense tracks correspond to HIPs, while faint tracks are produced by MIPs, since the aligned grains of these last tracks are separated. The local energy deposition (dE/dx) of each track can then be

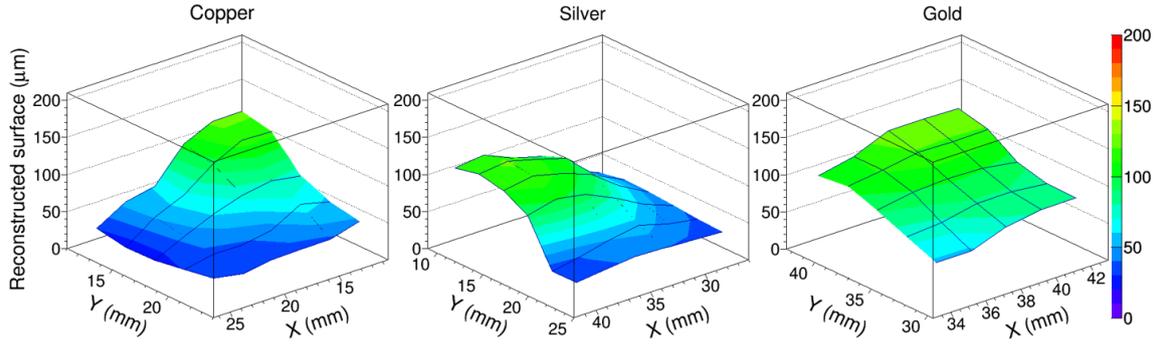


Figure 5. Surface topography for copper, silver and gold targets obtained from the reconstructed vertices. The vertical scale refers to the distance from the emulsion film.

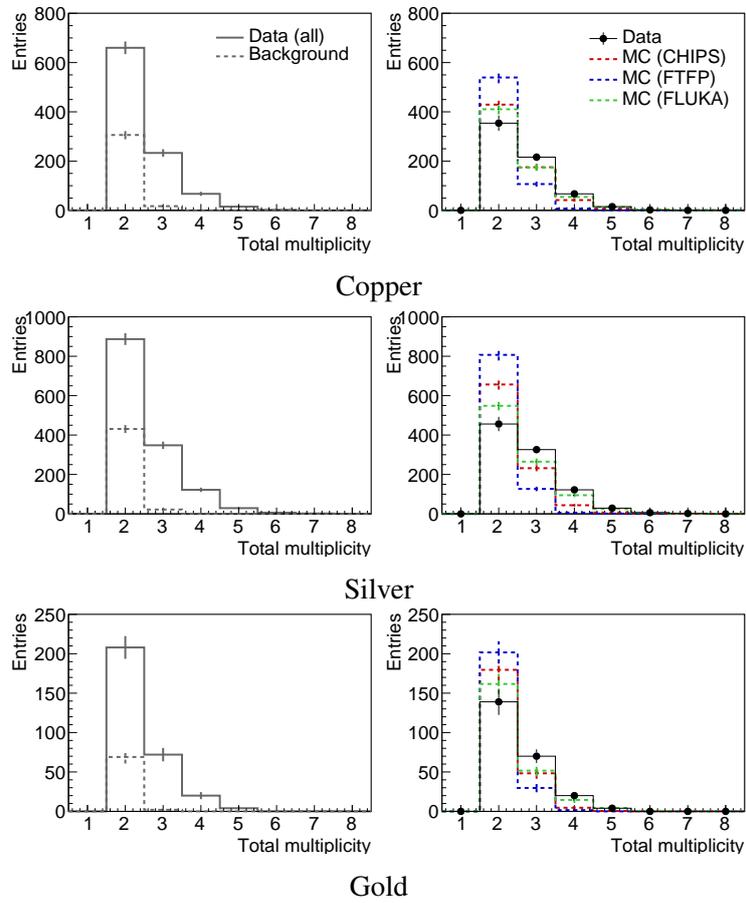


Figure 6. Left: Background contributions to the total multiplicities for the different targets. Right: Multiplicity distributions after subtraction of the background. The colored histograms show the Monte Carlo predictions.

assessed in terms of signal density (S.D.) along the reconstructed tracks, using

$$S.D. = \sum_{x,y,z \in C} S_{xyz}/L.$$

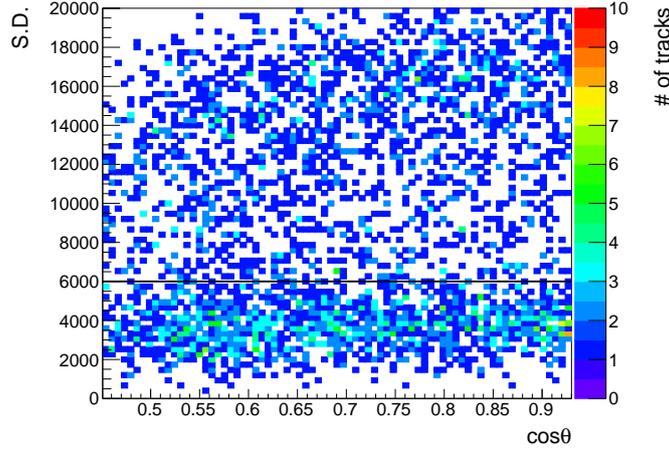


Figure 7. Signal density (S.D.) distribution as a function of track angle with respect to the beam direction. Tracks below the black line are defined as being minimally ionizing.

	Average multiplicity for MIPs				Average multiplicity for HIPs			
	Data	CHIPS	FTFP	FLUKA	Data	CHIPS	FTFP	FLUKA
Copper	1.07 ± 0.07	$0.98^{+0.07}_{-0.09}$	$1.59^{+0.09}_{-0.14}$	$0.83^{+0.08}_{-0.11}$	1.54 ± 0.07	$1.46^{+0.09}_{-0.07}$	$0.60^{+0.14}_{-0.10}$	$1.68^{+0.11}_{-0.08}$
Silver	1.02 ± 0.06	$1.04^{+0.08}_{-0.09}$	$1.64^{+0.09}_{-0.13}$	$0.73^{+0.07}_{-0.09}$	1.71 ± 0.07	$1.33^{+0.09}_{-0.08}$	$0.51^{+0.14}_{-0.09}$	$1.87^{+0.09}_{-0.07}$
Gold	0.92 ± 0.09	$1.21^{+0.10}_{-0.11}$	$1.75^{+0.09}_{-0.13}$	$0.81^{+0.07}_{-0.11}$	1.60 ± 0.09	$1.04^{+0.11}_{-0.09}$	$0.39^{+0.13}_{-0.09}$	$1.60^{+0.11}_{-0.06}$

Table 1. Measured average multiplicity and Monte Carlo predictions by CHIPS, FTFP (GEANT4) and FLUKA. Statistical errors and uncertainties in the background estimation are combined and reported for each set of data. The errors in the predictions account for uncertainties in dE/dx classification.

Here, x , y and z are the coordinates of voxels in the 3D image data. C is a group of voxels in a cylinder along the track, and S_{xyz} is an 8-bit grey-scale signal of the voxel. L is the length of a reconstructed track in the 3D image data. The S.D. is proportional to the dE/dx of the particle and does not depend on the angle. However, there is a saturation effect for higher values of dE/dx . Figure 7 shows that the S.D. distribution of tracks revealed a peak at 3000 for MIPs. As the simulated dE/dx distribution of MIPs peaked at $1.2 \text{ MeV}\cdot\text{g}\cdot\text{cm}^{-2}$, we define particles with dE/dx smaller than $2.4 \text{ MeV}\cdot\text{g}\cdot\text{cm}^{-2}$, corresponding to an S.D. below $6000 \mu\text{m}^{-1}$, as being MIPs. The complementary particles are defined as HIPs.

Figure 8 shows the track multiplicity distributions for MIPs and HIPs. The errors on the data are statistical. The histograms represent the Monte Carlo predictions by CHIPS, FTFP and FLUKA. The error bars of the Monte Carlo predictions account for uncertainties in the dE/dx classification. These uncertainties were estimated using simulations, which smeared the threshold for assigning tracks to either class of ionizing particles by 20% and checked for effects on the multiplicity distributions for MIPs and HIPs. The statistical errors for the simulations were 0.01-0.02, which are significantly smaller than the errors reported above. The average measured multiplicities are summarized in Table 1. Both CHIPS and FLUKA are in good agreement for copper, particularly in the case of MIPs. Neither CHIPS nor FTFP accurately describe particle multiplicities for

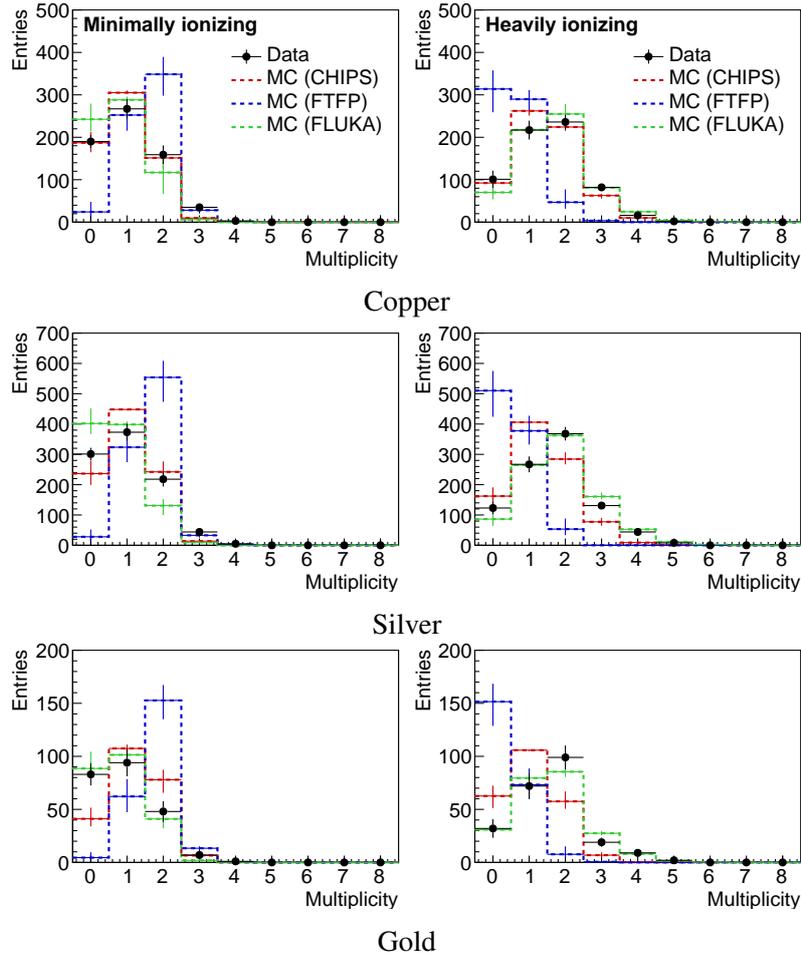


Figure 8. Reconstructed multiplicity distributions for annihilations in copper, silver and gold foils for MIPs (left) and HIPs (right). The colored histograms show the Monte Carlo predictions by CHIPS, FTFP and FLUKA. The error bars of the histograms account for uncertainties in the dE/dx classification.

annihilations on silver and gold nuclei, while FLUKA more closely reproduce the data than the other models.

The mean values of particle multiplicities measured for the three target materials are shown in Figure 9 as a function of atomic number along with the simulation outcome. Results obtained for MIPs with the FTFP model do not agree with our experimental data for any material, while those obtained with both CHIPS and FLUKA are in fair agreement as far as copper is concerned, although only FLUKA reproduces the higher atomic number behavior. Good agreement with CHIPS was also found for annihilation on bare emulsions and for aluminium [3]. Multiplicities related to HIPs are well described by the FLUKA simulation, while the CHIPS and FTFP models clearly underestimate the number of particles produced by antiproton annihilation.

4. Conclusions

The goal of the study presented in this paper was to measure the products of low-energy antiproton

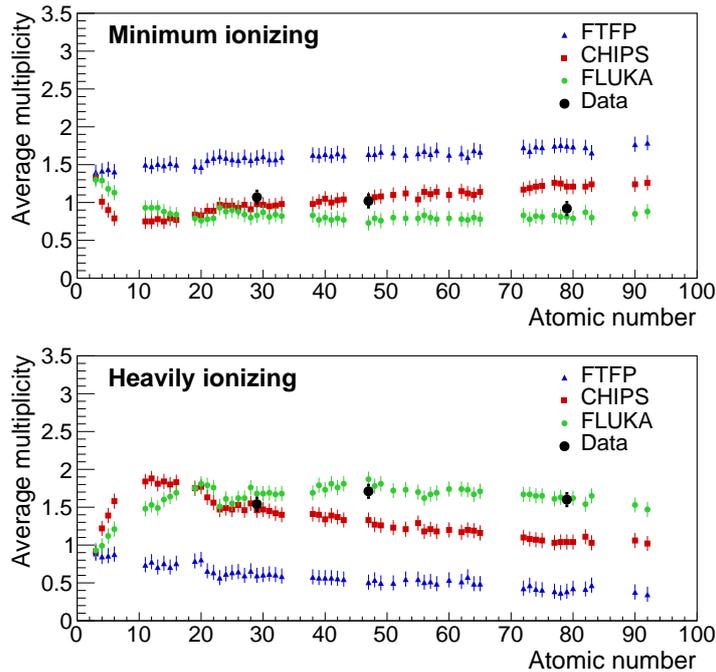


Figure 9. Particle multiplicity from antiproton annihilations as a function of atomic number for MIPs (top) and HIPs (bottom).

annihilation on different materials, utilizing a secondary beam line of CERN’s Antiproton Decelerator in the AEGIS experimental area. The characteristics (e.g. hadronization and fragmentation multiplicities) of low-energy antiprotons annihilating on nuclei are not well known, and experimental data are needed to validate models used by simulation packages such as GEANT4 and FLUKA. We exposed several thin targets (Cu, Ag and Au) to the antiproton beam and measured fragment tracks using emulsion detectors with a vertex position resolution at the level of a few micrometers, which allowed the separation between minimally and highly ionizing particles. The fragment multiplicities we measured were not well reproduced by the different models used in Monte Carlo simulation with the exception of FLUKA, which is in good agreement with the particle multiplicities for both minimally and heavily ionizing particles. Future measurements with more materials are needed to gain a better understanding of antinucleon annihilations also on low- Z materials, and to obtain a full description in terms of particle types, multiplicities, as well as energy, for a more complete benchmarking of Monte Carlo simulations.

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